## Stability of line solitons for KP-II

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The KP-II equation

$$\partial_x(\partial_t u + \partial_x^3 u + 3\partial_x(u^2)) + 3\sigma \partial_y^2 u = 0 \quad \text{for } t > 0 \text{ and } (x, y) \in \mathbb{R}^2$$
 (KP)

is a generalization to two spatial dimensions of the KdV equation

$$\partial_t u + \partial_x^3 u + 3\partial_x (u^2) = 0 \quad \text{for } t > 0 \text{ and } x \in \mathbb{R}.$$
 (1)

Obviously, any solutions of (1) are solutions of (KP) which are uniform in y. The KP equation takes slow variations of the transversal direction into account. If  $\sigma = 1$ , then (KP) is called KP-II and it describes the motion of shallow water waves with weak surface tension.

The KdV equation has a two parameter family of solitary wave solutions

$$\{\varphi_c(x-2ct-\gamma) \mid c>0, \gamma \in \mathbb{R}\}, \quad \varphi_c(x) = c \operatorname{sech}^2\left(\sqrt{\frac{c}{2}}x\right).$$

The KP-II equation is expected to explain stability of these solitary wave solutions to the transversal perturbations (see [7]). When  $\varphi_c(x - 2ct)$  is considered as a solution of the KP equations, we call it a *line soliton*. In this talk, I will discuss stability of line soliton solutions for the KP-II equation. The KP-II equation is known to be well-posed in  $L^2(\mathbb{R}^2)$  on the background of line solitons ([11]).

The KP equation (KP) is an integrable system as well as the KdV equation and have conserved quantities such as

$$\int_{\mathbb{R}^2} u(t, x, y)^2 \, dx \, dy \quad (\text{momentum}),$$
$$\frac{1}{2} \int \left( u_x^2(t, x, y) - 3\sigma (\partial_x^{-1} \partial_y u(t, x, y))^2 - 2u^3(t, x, y) \right) \, dx \, dy \quad (\text{Hamiltonian})$$

For the KP-I equation, that is (KP) with  $\sigma = -1$ , the first two terms of the Hamiltonian have the same sign and it was shown by [6, 14] that there exists a stable ground state for the KP-I equation, whereas the KP-II equation has no traveling wave solutions that belong to  $L^2(\mathbb{R}^2)$ ([5]).

On the other hand, it is known that the line solitons of the KP-I equation are unstable ([16]) and line solitons of the KP-II equation are linearly stable. Because the quadratic part of the Hamiltonian of the KP-II equation, which is the dominant part, is indefinite, it is more natural to explain stability of line solitons by using propagation estimates such as [13] rather than by using variational arguments such as [3].

The main difference between stability analysis of the KdV 1-soliton or the stability of line soliton with the y-periodic boundary condition is that  $\varphi_c(x-2ct)$  does not have the finite  $L^2$ -mass because it is not localized in the y-direction. As a consequence, the linearized operator of the KP-II equation around the line soliton has a family of continuous eigenvalues converging to 0 (see e.g. [1, 2]) in exponentially weighted space, whereas 0 is an isolated eigenvalue of the linearized KdV operator around a solitary wave in exponentially weighted space.

Since we have continuous spectrum converging to 0, the modulations of the speed and the phase shift of line solitons for the KP-II equation cannot be described by ODEs as KdV ([13]) or the KP-II equation posed on  $\mathbb{R}_x \times \mathbb{R}_y/(2\pi\mathbb{Z})$  ([10]).

I find that for the KP-II equation posed on  $\mathbb{R}^2$ , modulation of the speed parameter c(t, y) and the phase shift x(t, y) cannot be uniform in y and their long time behavior is described by the Burgers equation. Note that similar modulation equations have formally derived for a two spatial dimensional Boussinesq model ([12]).

My results are the following ([9]).

**Theorem 1** Let  $c_0 > 0$  and  $a \in (0, \sqrt{c_0/2})$ . Then there exist positive constants  $\varepsilon_0$  and C satisfying the following: if  $u(0, x, y) = \varphi_{c_0}(x) + v_0(x, y), v_0 \in H^1(\mathbb{R}^2)$  and  $\varepsilon := \|e^{ax}v_0\|_{L^2(\mathbb{R}^2)} + \|e^{ax}v_0\|_{L^1_yL^2_x} + \|v_0\|_{L^2(\mathbb{R}^2)} < \varepsilon$   $\varepsilon_0$ , then there exist C<sup>1</sup>-functions c(t, y) and x(t, y) such that for  $t \ge 0$ ,

$$\|u(t,x,y) - \varphi_{c(t,y)}(x - x(t,y))\|_{L^2(\mathbb{R}^2)} \le C\varepsilon, \qquad (2)$$

$$\sup_{y \in \mathbb{R}} (|c(t,y) - c_0| + |x_y(t,y)|) \le C\varepsilon (1+t)^{-1/2},$$
(3)

$$\|x_t(t,\cdot) - 2c(t,\cdot)\|_{L^2} \le C\varepsilon(1+t)^{-3/4},$$
(4)

$$\left\| e^{ax}(u(t, x + x(t, y), y) - \varphi_{c(t,y)}(x)) \right\|_{L^2} \le C\varepsilon (1+t)^{-3/4}.$$
 (5)

We find that c(t, y) and  $\partial_y x(t, y)$  behave like a self-similar solution of the Burgers equation around  $y = \pm \sqrt{8c_0}t$ .

**Theorem 2** Let  $c_0 = 2$  and let  $v_0$  and  $\varepsilon$  be the same as in Theorem 1. Then for any R > 0,

$$\left\| \begin{pmatrix} c(t, \cdot) \\ x_y(t, \cdot) \end{pmatrix} - \begin{pmatrix} 2 & 2 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} u_B^+(t, y+4t) \\ u_B^-(t, y-4t) \end{pmatrix} \right\|_{L^2(|y\pm 4t| \le R\sqrt{t})} = o(t^{-1/4})$$

as  $t \to \infty$ , where  $u_B^{\pm}$  are self similar solutions of the Burgers equation

$$\partial_t u = 2\partial_y^2 u \pm 4\partial_y(u^2)$$

such that

$$u_B^{\pm}(t,y) = \frac{\pm m_{\pm} H_{2t}(y)}{2\left(1 + m_{\pm} \int_0^y H_{2t}(y_1) \, dy_1\right)}, \quad H_t(y) = (4\pi t)^{-1/2} e^{-y^2/4t},$$

and that  $m_{\pm}$  are constants satisfying

$$\int_{\mathbb{R}} u_B^{\pm}(t, y) \, dy = \frac{1}{4} \int_{\mathbb{R}} c(0, y) \, dy + O(\varepsilon^2) \, .$$

The KP-II equation (KP) is invariant under a change of variables

$$x \mapsto x + ky - 3k^2t + \gamma$$
 and  $y \mapsto y - 6kt$  for any  $k, \gamma \in \mathbb{R}$ , (6)

and has a 3-parameter family of line soliton solutions

$$\mathcal{A} = \{\varphi_c(x + ky - (2c + 3k^2)t + \gamma) \mid c > 0, k, \gamma \in \mathbb{R}\}$$

Thanks to propagations of the local phase shifts along the crest of line solitons, the set  $\mathcal{A}$  is not stable in  $L^2(\mathbb{R}^2)$ .

**Theorem 3** Let  $c_0 > 0$ . There exists a positive constant C such that for any  $\varepsilon > 0$ , there exists a solution of (KP) such that  $||u(0, x, y) - \varphi_{c_0}(x)||_{L^2} < \varepsilon$  and

$$\liminf_{t \to \infty} t^{-1/4} \inf_{v \in \mathcal{A}} \|u(t, x, y) - v\|_{L^2(\mathbb{R}^2)} \ge C\varepsilon.$$

I remark that similar problems have already been studied for nonlinear heat equations ([15, 8]) and also for kink solutions to a wave equation ([4]). However, we need some correction of the phase shift for the KP-II equation, which does not appear in these former results because the transversal dimension and the power of nonlinearity is lower in our problem.

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